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AFML-TR-67-143 PART III

## **ELASTIC-PLASTIC FRACTURE MECHANICS**

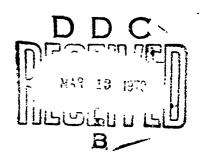
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Battelle Memorial Institute

Columbus Laboratories

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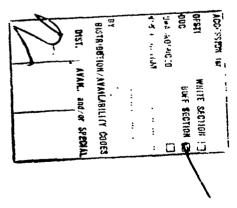
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## **ELASTIC-PLASTIC FRACTURE MECHANICS**

G. T. HAHN, A. R. ROSENFIELD, and M. SARRATE®

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 $<sup>{}^{</sup>ullet} M.$  Sarrate was on leave of absence from the Argentine Atomic Energy Commission.

### **FOREWORD**

This report was prepared by personnel of the Columbus Laboratories of Battelie Memorial Institute under USAF Contract No. AF 33(615)-3565. The contract was initiated under Project No. 7351, "Metallic Materials", Task No. 735106, "Behavior of Metals". The contract was administered under the direction of the Metals and Ceramics Division of the Air Force Materials Laboratory, Wright-Patterson Air Force Base, with Mr. A. W. Brisbane acting as project engineer.

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The work reported began on 1 April 1968, and was concluded on 30 September 1969. This report was submitted by the authors 30 September, 1969.

This technical report has been reviewed and is approved.

W. J. TRAPP

Chief, Strength and Dynamics Branch Metals and Ceramics Division Air Force Materials Laboratory

### ABSTRACT

This report describes experiments that delineate the plastic zone of a growing fatigue crack and its relation to the zone of a monotonically loaded, stationary crack. A compilation of cyclic crack growth measurements is also presented. Preliminary results of the zone studies and the data compilation are examined with respect to the mechanism of cyclic growth and the contribution of metallurgical factors.

Two techniques, etch pitting and interferometry, are used to reveal the plastic zones produced by growing fatigue cracks in Fe-3Si Steel (under plane strain) and in cold worked steel (under plane stress), respectively. In addition, a method of simulating cyclic crack growth under controlled conditions by cutting-in slits under load with intermittent load reversals is demonstrated. The preliminary results indicate that the plastic deformation generated by each loading cycle is similar to the deformation of the monotonic zone.

The results provide a basis for extending numerical treatments of the monotonically loaded crack to the fatigue crack problem. Estimates are made in this way of the number of plastic strain cycles experienced by the material in front of a growing fatigue crack and also the crack-tip displacement. These values together with the compilation of crack growth measurements are examined in the light of two fatigue crack growth mechanisms: (1) irreversible crack blunting, and (2) Manson-Coffin type strain accumulation. This shows that both mechanisms can account for the value of the stress intensity exponent, m  $\approx$  2 observed in Regime No. 1 (the high cycle-low stress portion of the crack growth spectrum). While neither mechanism easily accounts for the invariance of crack growth rates in Regime No. 2, the existing observations are more easily rationalized in terms of blunting. A possible explanation for the greater stress sensitivity and higher growth rates in Regime No. 2 is offered in terms of an increasing length of active crack front.

Both the theoretical treatment and the bulk of the test data show that the resistance to fatigue crack growth in Regime No. 1 is relatively insensitive to composition and metallurgical structure, provided environmental effects are not involved. In contrast, significant improvements in fatigue life could be achieved simply by postponing the onset of Regime No. 2, a feature that should depend on metallurgical variables.

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## TABLE OF CONTENTS

	Page
ı.	EXPERIMENTS ON THE NATURE OF THE FATIGUE CRACK PLASTIC ZONE 1
	INTRODUCTION
	EXPERIMENTAL PROCEDURE
	Plane Stress
	RESULTS
	Plane Stress
	DISCUSSION
	CONCLUSIONS
	REFERENCES
II.	APPENDIX ACompilation of Cyclic Crack-Growth Rate Measurements 45
	APPENDIX BSimplified Formulations of the Cyclic Crack Growth Mechanisms
	APPENDIX CAugmented Dugdale Model
	LIST OF ILLUSTRATIONS
Figu	<u>ce</u>
1	Cantilever Fracture Toughness Specimen
2	Stress-Strain Characteristics of the Fe-3Si Steel at 100 C 3
3	Strain Distribution Ahead of a Slit Revealed Under Load by Interferometric Fringe Pattern
4	Influence of Cyclic Loading, $\frac{\Delta K}{\sigma_Y} = 0.225 \sqrt{\text{in}}$ , ON THE STEEL FOIL 7
5	Influence of Cyclic Loading on the Crack Tip $\epsilon_z$ -STRAIN PROFILES 8
6	Plastic Zone Produced in the Interior of Specimen 3P-20 By the Fatigue Simulation Procedure

# LIST OF ILLUSTRATIONS (Continued)

		Page
Figure		
7	Plastic Zone Produced in the Interior of Specimen 3P-19 By Monotonic Loading Simulation	13
8	Plastic Zone Produced in the Interior of Specimen 3P-18 By the Stable Growth Simulation Procedure	15
9	Plastic Zone of a Mixed Fatigue-Cleavage Crack Produced In Specimen 3P-24 By 21,770 Cycles with a Stress Intensity Range $\frac{\Delta K}{\sigma_{_{\mathbf{V}}}} = 0.7\sqrt{\text{in}}$	18
10	Plastic Zone of a Mixed Fatigue-Cleavage Crack Produced in Specimen 3P-25 By 9930 Cycles with a Stress Intensity	
	Range $\frac{\Delta K}{\sigma_{Y}} = 0.7\sqrt{\text{in}}$	24
11	Comparison of Fatigue Crack- and Stationary, Virgin Crack-Near-Tip Zones	28
12	Schematic of Monotonically Loaded Crack Tip Plastic Zone	31
13	Simplified Schematic Representation of the Origin of Regimes No. 1 and No. 2	38
A-1	Schematic of Stress-Intensity-Growth Rate Spectrum with 2 Regimes and Definitions of Terms in Table A-1	48
C-1	Augmented Dugdale Crack Model	51
	LIST OF TABLES	
<u>Table</u>		
1	Summary of COD and Crack-Tip Strain Values	9 ·
2	Summary of Plastic Zone Size-Values Obtained from Etching Experiments	17
3	Estimates of $N$ and $N$ , the Numbers of Plastic Strain Cycles and $\beta$ , the Ratio of Crack Growth Per Cycle to Crack-Opening Displacement	32
4	Summary of Selected Cyclic Crack Growth Resistance Values	35
A-1	Summary of Cyclic Crack Growth Data	45

# EXPERIMENTS ON THE NATURE OF THE FATIGUE CRACK PLASTIC ZONE

### INTRODUCTION

Until recently, few generalizations could be made about fatigue crack growth. Although much of the data has been expressed as a power relation:  $\frac{da}{dt} = A(\Delta K)^m$ , neither of the curve-fitting parameters is, in general, a constant. Both m, and A depend on the stress range, the material, the test environment, and the exact details of the testing procedure. In the absence of correlations, it has been difficult to identify the metallurgical influences on the fatigue crack growth resistance.

Now, with the accumulation of more test data, (1-27) patterns are beginning to emerge. The compilation in Appendix A contains evidence of two growth-rate regimes within which many alloys display striking similarities. Progress is also being made on the analytical front; see Rice<sup>(28)</sup> for a comprehensive review and more recent treatments by Rice and Johnson<sup>(29)</sup>, Erdogan<sup>(30)</sup>, and Liu and coworkers<sup>(22,31)</sup>. These analyses hinge on the two mechanisms of crack advance that have been proposed: (1) irreversible plastic blunting<sup>(32-34)</sup> and (2) Manson<sup>(35)</sup>-Coffin<sup>(36)</sup>-type damage accumulation. They also depend on descriptions of the strain distribution in the locale of the fatigue crack<sup>(34,37,38)</sup>.

The experimental work described in this report deals with the plastic zone of a growing fatigue crack and its relation to the zone of a monotonically loaded, stationary crack. This subject is virtually unexplored because of the experimental difficulties, and only one set of measurements, by Liu and Iino (22) are known to the authors. The present study examines ways of applying two techniques: etch pitting and interferometry, to reveal the plastic zones produced by fatigue cracks under plane strain and plane stress conditions. Preliminary results are reported and these indicate that the plastic deformation generated by each loading cycle is similar to the zone of a stationary crack loaded monotonically.

On this basis, theoretical treatments of the wonotonically loaded crack are tentatively extended to the fatigue crack problem. Simplified formulations of the plastic blunting and damage accumulation are obtained in this way. The efficiency of the blunting mechanism and the number of plastic strain cycles experienced by material in front of the crack is estimated. This shows that both mechanisms can account for the value of the stress intensity exponent,  $m\approx 2$  observed in Regime No. 1 (the high cycle-low stress portion of the crack growth spectrum). While neither mechanism easily accounts for the invariance of crack growth rates in Regime No. 1, the existing observations are more easily ration-

 $t = \frac{da}{dN}$  is the cyclic crack growth rate,  $\Delta K$  is the cyclic stress intensity variation, A is a materials parameter and m is a numerical constant.

alized in terms of blunting. A possible explanation is offered for the higher m-values and growth rates of Regime No. 2 (the low cycle-high stress range). Implications with respect to the metallurgical origins of the cyclic crack growth resistance and the prospects of improving it are discussed.

### EXPERIMENTAL PROCEDURE

### Plane Stress

Plastic zones under plane stress were revealed in an 0.00175 in. thick, cold rolled, mild steel foil—yield strength: 105,000 psi and toughness:  $K_c = 34 \text{ ksi } \sqrt{\text{in}}$ —whose fracture behavior has been reported on before (39-40). The behavior of the foil approximates two idealizations: elastic-perfectly plastic behavior and plane stress deformation. Because plane stress conditions are approached, the deformation within the plastic zone ahead of a crack is mainly transverse to the surface and easily measured.

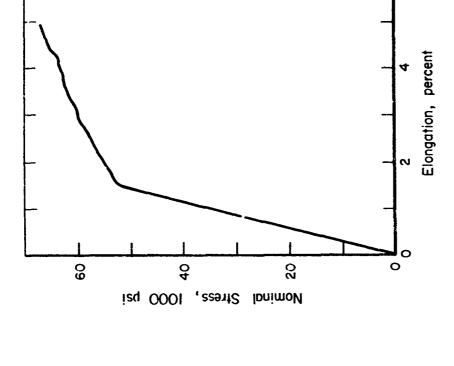
The foil specimen was 4.0" long by 3.9" wide with a centrally located notch 0.220" long cut in by spark machining. It was prepolished to give a smooth surface and mounted in a small tensioning device which could be placed in the stage of an interference microscope. The strain distribution could thus be detected under load as well as after unloading. The specimen was loaded by means of a screw-and-lever arrangement, the load being monitored with an electric resistance strain gage (more detailed descriptions of the experimental methods are given in References (39) and (40)). The fatigue loading corresponds to cycling between stress-intensity levels of  $K/\sigma_V = 0.025-0.25\sqrt{in}$ . After varying numbers of cycles, the loading was interrupted and interference patterns of the strain distribution at maximum and minimum loads were made.

### Plane Strain

Studies of plastic zones under plane strain employed 1/2 in.-thick, Fe-3Si steel plates as the model material<sup>†</sup>. Descriptions of the material and the procedure for revealing zones on the surface and on interior sections by etching, as well as the zones displayed by virgin, stationary cracks in this steel after one cycle load-unload are given in Reference (41). Unlike the earlier work, the present studies were conducted with the DCB (double cartilever beam) specimen shown in Figure 1. Stress intensity values were calculated from the Mostovoy relation<sup>(42)</sup> and checked against compliance measurements. In an effort to avoid cleavage, the tests were carried out at 100 C, but this strategy was only partially successful. Cleavage can probably be avoided in future experiments by lowering the silicon content. The tensile properties of the steel at 100 C, in Figure 2, show that the yield stress is 54,000 psi at this temperature.

In the first experiment fatigue crack growth was simulated by cutting in a sharp slit under load with intermittent unloading. The specimen was first

<sup>†</sup> Nominal Composition Si: 3.4%, C: 9.02% remainder Fe.





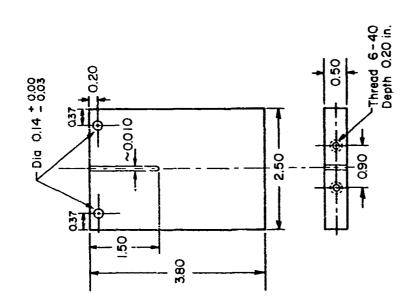


FIGURE 1. CANTILEVER FRACTURE TOUGHNESS SPECIMEN

loaded to a predetermined stress intensity level,  $\frac{K}{\sigma_V} = 0.7 \, \text{Vint}$  in a horizontal tensile testing machine. Then the slit was extended 0.0008 in. by spark machining to under load. At this point, the specimen was unloaded, and reloaded to  $\frac{K}{\sigma_v} = 0.7\sqrt{1n}$ , and the sequence of steps, cutting, unloading, and reloading 50 times in quick succession. Effects of the stress veversals within zone and the plastic region left behind the growing slit were discerned by comparisons with the zones produced by:

- (i) A virgin slit subjected to one cycle of loading, and by
- (ii) A slit cut in under load with the load controlled to produce a constant stress intensity at the crack tip which are described more fully elsewhere. (43)

In the second type of experiment, actual fatigue cracks were generated in pre-slitted DCB specimens and grown by cycling between a small preload and a maximum load and deflection controlled to give a fixed stress intensity range of  $\frac{\Delta K}{\sigma_V}$  = 0.7  $\sqrt{\text{in}}$ . These experiments were performed at 10 cps in a standard hydraulic fatigue testing machine. The zones associated with these cracks were then revealed by etching.

### RESULTS

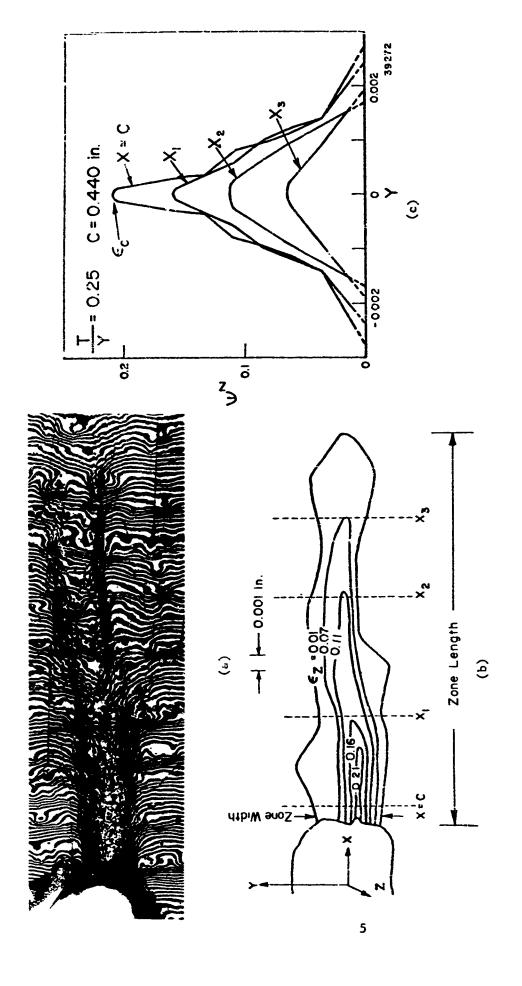
### Plane Stress

Figure 3 illustrates the condition of a sample under load after the first 1/2 fatigue cycle and the information derived from the interferometric fringe pattern (Figure 3a) it. No crack growth was discernible in the steel foil after

<sup>†</sup> K is the stress intensity, and  $\sigma_{\rm Y}$  is the yield stress.

<sup>††</sup> The slits were spark machined with a 0.005 in. diameter copper wire which produces a slit  $\sim 0.010$  in. wide with a  $\sim 0.005$  in. root radius.

The patt on reflects the distribution of the through-the-thickness strain €2, and this can be converted into an isostrain contour map (Figure 3b) which shows the general shape of the zone, the zone length and the zone width at the crack (or slit) tip, as well as the distribution of strain within. Since the patterns are difficult to interpret right at the crack (or slit) tip, values attributed to the crack tip, i.e. at x = c, were always measured 0.001 in. from the tip. The  $\epsilon_z$ -profile at a particular distance from the crack tip can also be constructed (Figure 3c). The peak strain displayed by the profile at the crack tip, x = c, is referred to here as the crack-tip strain or  $\epsilon_c$ . The area within this profile, i.e.,  $-2\frac{1}{2}$   $\epsilon_z$  dY represents the COD, or crack opening displacement. The fact that the zones are necked largely precludes in-plane strain:  $\epsilon_{x} = 0$ . This and the constancy of volume give the result:  $\epsilon_{y} = -\epsilon_{z}$  (see Figure 4b for coordinates). Consequently, the COD (the total Y-displacement at the crack tip arising from plastic deformation) is:  $COD = 2\sqrt{\epsilon_Y d_Y} = -2\sqrt{\epsilon_Z} dY$ . Note that the zone in Figure 3 represents the deformation produced by the first half of a single load-unload cycle involving a stress intensity range  $\frac{\Delta K}{C_Y} = 0.3\sqrt{in}$ , the zone corresponding to the first half of the cycle employed in the fatigue experiments  $\frac{\Delta K}{C_Y} = 0.25\sqrt{in}$  was similar in character but about 15% smaller.



pattern displayed by sample with c = 0.440 in, loaded to T/Y = 0.25, K/Y = 0.30 Vin., (b) the isostrain contours derived from the fringes, and (c) the  $\epsilon_z$ -strain profile at four points along the plastic zone. (a) and (b) 150% STRAIN DISTRIBUTION AHEAD OF A SLIT REVEALED UNDER LOAD BY INTERFEROMETRIC FRINGE PATTERN: FIGURE 3.

the ten cycles of loading, but when the fatiguing was interrupted for zone measurement after twenty cycles the crack had extended slightly at either end (Figure 4a). Eventually, the growth rate settled down to a roughly constant value of  $\sim 2 \times 10^{-4}$  in/cycle.†† This rate is not unlike those reported for high strength structural alloys at comparable  $\frac{\Delta K}{CY}$  - levels. It also appears that the initial growth rate is slightly larger than during steady-state growth.

Qualitatively, the plastic zone had the same long, narrow shape after many cycles of loading as it did after the first cycle before crack growth was observed. Both the zone length and the zone width increase as a result of cyclic loading. The zone length after one cycle is within the range (but somewhat on the high side) of our previous measurements. During the first few cycles of loading, the length remains roughly constant and only appears to increase upon the onset of crack growth (Figure 4b). During crack growth, the length again remains relatively constant at a value  $\sim 60\%$  greater than the original value. On the other hand, the zone appears to broaden even before crack growth starts. However, the width also attains a saturation value which is  $\sim 60\%$  greater than after the initial cycle (Figure 4c).

The on-load and off-load strain profiles of the first fatigue cycle, including the COD and  $\epsilon_{\rm C}$ , appear to be identical within the experimental error (see Figure 5a). This result as well as the disparate behavior of the two ends of the crack are puzzling. While the crack tip strain profile on the right end assumed the steady state value prior to the fifth cycle (Figures 5b and 5d), the profiles on the left end appeared to be unaffected up to the onset of crack extension sometime between the 10th and 20th cycle. More experiments are needed to establish the significance of these results which may have been affected by heterogeneities in the foil, asymmetry of loading, and possibly, some tendency for the foil to buckle.

The results obtained after the onset of cracking appear to be more consistent. Figures 5b and d illustrate that a steady state was achieved quickly. Steady state values of the off-load COD and  $\epsilon_{\rm C}$  (which remained essentially unchanged throughout) were  $\sim$  60% of the on-load values (see Figure 5c and Table 1). The on-load values were about 50% larger than those attending crack extension and stable crack growth during monotonic loading. Both the larger zones and the enhanced ductility may be connected with the cyclic softening (33) of the cold worked steel foil, an effect that has been observed on other materials. While these observations also lend support to the idea that fatigue crack growth proceeds with a constant of critical COD, they also suggest that this COD need not be the same as the value associated with the ordinary crack extension.

In the present experiments the material directly in front of the crack experienced about 200 cycles of reverse plastic strain.† Since the maximum strain range per cycle =2 · [ $\epsilon_{\rm c}$ (on-load)- $\epsilon_{\rm c}$ (off-load)] = 0.25 at the crack tip (see Table 1), the average plastic strain range over the entire zone length

<sup>†</sup> The number of plastic strain cycles  $N = \rho \left(\frac{da}{dN}\right)^{-1}$ , where  $\rho$  is the plastic zone length and  $\frac{da}{dN}$  the crack advance per cycle. At steady state,  $\rho = 0.040$  in., and  $\frac{da}{dN} = 2.10-4$  in.

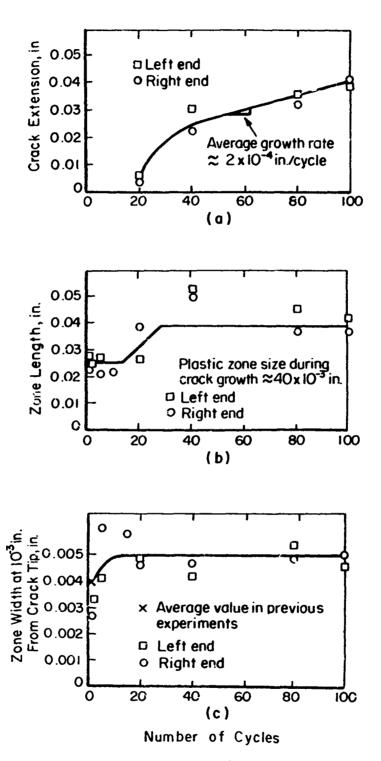


FIGURE 4. INFLUENCE OF CYCLIC LOADING,  $\frac{\Delta K}{\sigma y}$  = 0.225  $\sqrt{\text{in}}$ , ON THE STEEL FOIL: (a) crack length, (b) plastic zone length, and (c) plastic zone width

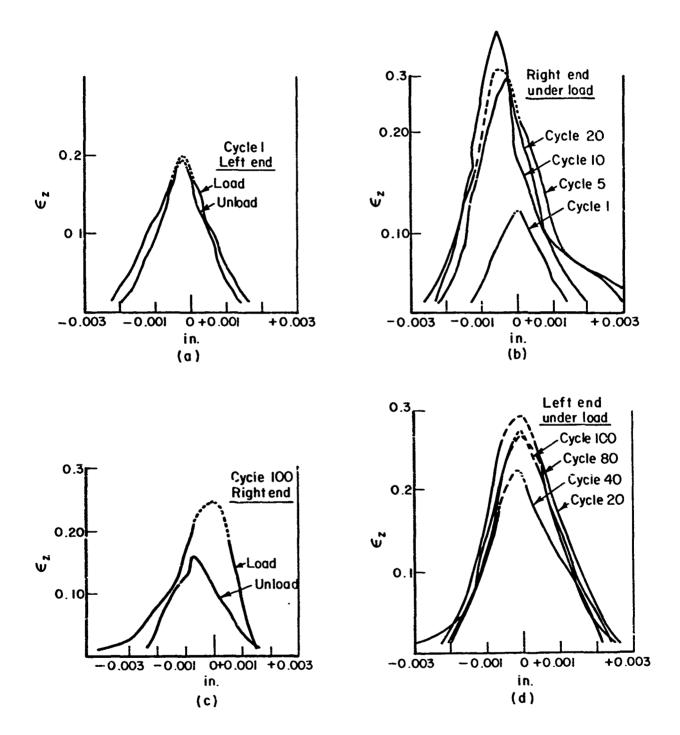


FIGURE 5. INFLUENCE OF CYCLIC LOADING ON THE CRACK TIP  $\epsilon$  -STRAIN PROFILES: (a) and (d) for left end, (b) and (c) for right end

TABLE 1. SUMMARY OF COD AND CRACK-TIP STRAIN VALUES

Type of Cracking	$\frac{K}{\sigma_{Y}}$ (in <sup>1/2</sup> )	COD (10 <sup>-4</sup> in)	€ <sub>c</sub>
Stable crack growth under monotonic loading (40)	0.32-0.35	2.6	0.24
Fatigue Crack Growth			
lst Cycle, Load-on	0.25	2.0	0.20
lst Cycle, Load-off	0.025	1.7	0.17
20th-100th Cycle, Load-on	0.25	3.9	0.31
20th-100th Cycle, Load-off	0.025	2.3	0.185

is about 1/2 this value, or  $\overline{\Delta \varepsilon}_{\rm p} \sim$  0.1, since the strain gradient is nearly linear (40). Thus, the total plastic strain accumulation prior to cracking,  $\overline{\Delta \varepsilon}_{\rm p} n \sim$  20 (85 times the value, 0.24, associated with failure under monotonic loading). Such strains could be accumulated if the material in front of the crack in the steel foil obeys a Manson-Coffin-type fatigue damage law (35-36). Furthermore, a strain range  $\overline{\Delta \varepsilon}_{\rm p} \sim$  0.1 is roughly consistent with expectations for a life of  $\sim$  100 cycles based on unnotched tests on a large number of materials (44).

### Plane Strain

The plastic zone produced in the DCB-specimen by the fatigue simulation procedure is shown in Figure 6. It can be compared with the zone of a stationary slit after one load-unload cycle reproduced in Figure 7 and with the zone of a slit cut-in under load at constant K without intermittent cycling in Figure 8. The size of the zone at the crack tip is essentially the same in all three cases (see Table 2)†. On the basis of the relation between zone size and COD shown to exist for simplified elastic plastic models†, the COD's and near tip strain distributions should also be similar. This is in accord with the close-ups of the etched tip of the stationary slit and the one cut in at constant-K.††† However, when the slit (or crack) is subjected to cyclic loading, the deformation from a number of zones will be superimposed at the crack tip (depending on the growth per cycle relative to the forward extent of the zone). The zone produced by cutting-in with intermittent unloading (Figure 6b) does, in fact, show a larger light etching zone at the crack tip††† consistent with this picture.

The plastic zones and fracture morphologies of cracks produced by actual cyclic growth are shown in Figures 9-11. It is apparent from the pictures of the plate midsections, Figure 9, that the cyclic growth process was accompanied by some cleavage cracking. The fatigue portions of the failures are characterized by extensive light-etching regions on either side of the crack reflecting the superposition of deformations and resulting strain accumulation, and by the very ragged, irregular crack path with numerous instances of branching (see Figure 9c,d,g, and 10c). In some places the photos show irregularities on the crack path on a scale approaching the crack advance per cycle: ~2.10<sup>-5</sup> in.

<sup>†</sup> For convenience, the quantity  $\rho$ , which is referred to as the zone size is defined here as the furthest extent of the etched region from the slit on the plate midsection measured normal to the plane of the slit.

<sup>††</sup>  $COD = \left(\frac{8}{7} \frac{\sigma_Y}{E}\right) \rho$ , where  $\rho$  is the zone size, E is the elastic modulus and  $\sigma_Y$  is the yield stress<sup>(43)</sup>. This relation is valid at relatively low nominal stress levels  $\sigma$ :  $\sigma < 0.7 \sigma_Y$ .

then the Fe-3Si Steel is strained more than  $\sim$  5%, it no longer etches dark. The interface between the dark and light etching region at the slit tip can be regarded as an isostrain contour corresponding to about 5%.

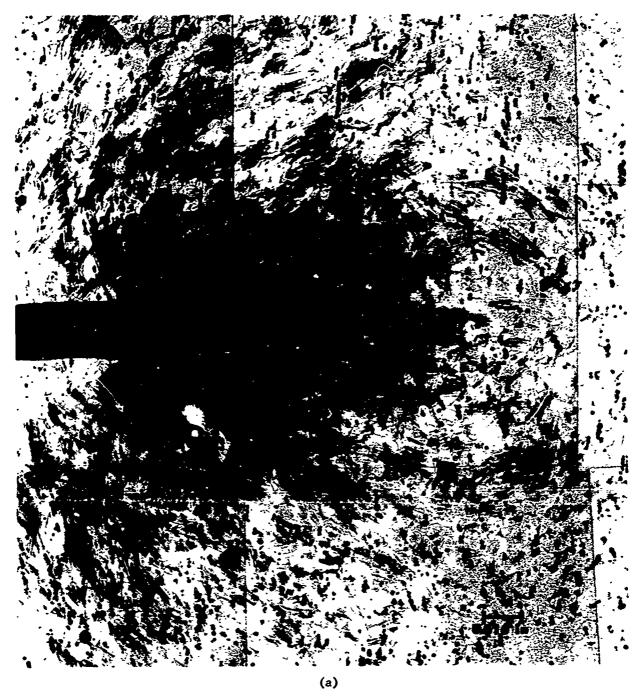


FIGURE 6. PLASTIC ZONE PRODUCED IN THE INTERIOR OF SPECIMEN 3P-20 BY THE FATIGUE SIMULATION PROCEDURE. The slit was cut in under loads corresponding to  $\frac{K}{\sigma_Y} = 0.7 \sqrt{\text{in.}} \text{ in } 50\text{--}0.0008 \text{ in.} \text{ increments with intermediate unloading:}$ 

(a) zone revealed on specimen midsection, and (b) close-up of slit tip



FIGURE 6. PLASTIC ZONE PRODUCED BY FATIGUE SIMULATION PROCEDURE (Continued)

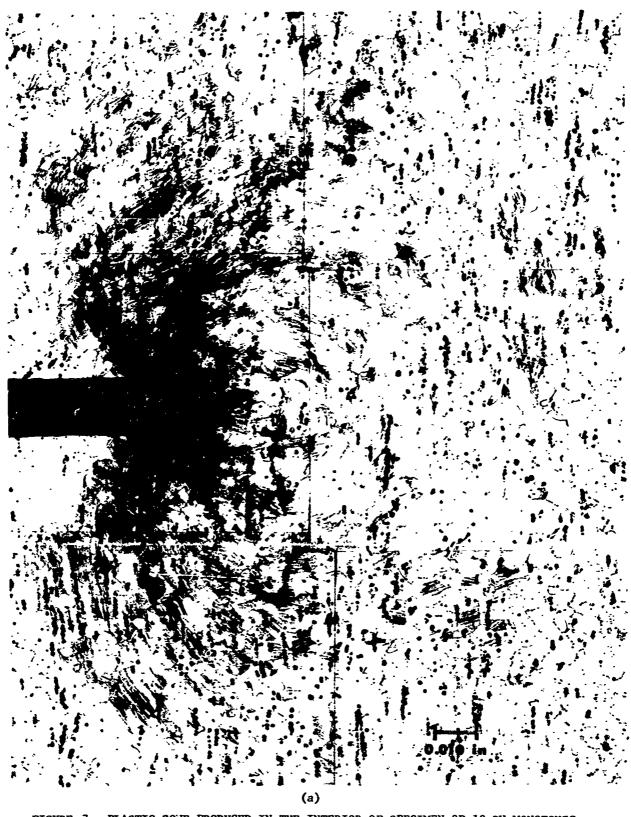


FIGURE 7. PLASTIC ZONE PRODUCED IN THE INTERIOR OF SPECIMEN 3P-19 BY MONOTONIC LOADING SIMULATION. Pre-cut, virgin slit was loaded once to  $\frac{K}{\sigma_V}=0.7\,\sqrt{in}$  and unloaded: (a) zone revealed on specimen midsection, and (b) close-up of slit tip



FIGURE 7. PLASTIC ZONE PRODUCED IN THE INTERIOR OF SPECIMEN 3P-19 BY MONOTONIC LOADING SIMULATION (Continued)



FIGURE 8. PLASTIC ZONE PRODUCED IN THE INTERIOR OF SPECIMEN 3P-18 BY THE STABLE GROWTH SIMULATION PROCEDURE. The slit was cut in under loads corresponding to a constant stress intensity,  $\frac{1}{C_Y} = 0.7 \sqrt{\text{in}}$ : (a) zone revealed on specimen midsection and (b) close-up of slit tip



PLASTIC ZONE PRODUCED IN THE INTERIOR OF SPECIMEN 3P-18 BY THE STABLE GROWTH SIMULATION PRO-CEDURE(43) (Continued) FIGURE 8.

TABLE 2. SUMMARY OF PLASTIC ZONE SIZE VALUES OBTAINED FROM ETCHING EXPERIMENTS

Spec. No.	Description of	$\frac{K_{-}}{(4n^{1/2})}$	Cyclic Growth Rate	Zone Si	Zone Size, P (in)
	Experiment	$\sigma_{\mathbf{Y}}$	(in)	Plate Surface	Plate Midsection
32-19	Pre-cut slit subjected to one load-unload cycle	7.0	1	0.058	0.075
3P-23	<del>.</del>	0.7	1	0.053	690.0
3P-18	Slit cut in under load at constant stress intensity	0.7	1	0.055	0.074
38-20	Slit cut in under load with intermediate load cycling	0.7	8 • 10-4	0.058	0.070
3P-24	Fatigue Crack Growth	7.0	~ 2 · 10-5	0.065	0.070
3P-25	=	0.7	1.2 · 10-5	970.0	090°0



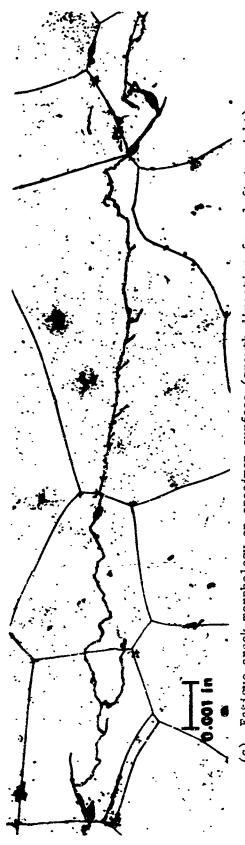
PLASTIC ZONE OF A MIXED FATIGUE-CLEAVAGE CRACK PRODUCED IN SPECIMEN 3P-24 BY 21,770 CYCLES WITH A STRESS INTENSITY RANGE  $\frac{\Delta K}{\sigma_{\mathbf{Y}}}=$  0.7  $\sqrt{1n}$  (CLEAVAGE PORTIONS ARE LABELED C)

FIGURE

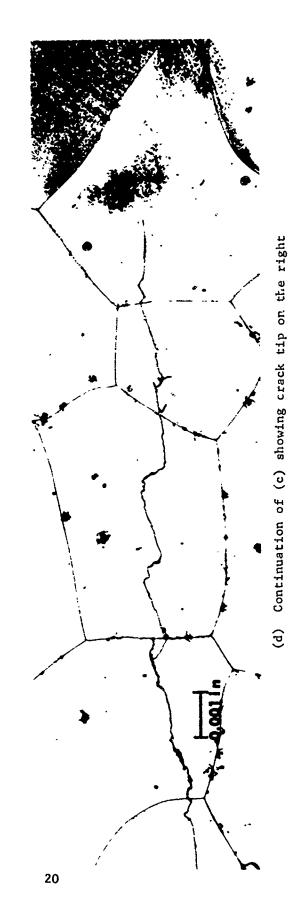


(b) Continuation of zone on specimen surface

FIGURE 9. PLASTIC ZONE OF A MIXED FATIGUE-CLEAVAGE CRACK IN SPECIMEN 3P-24 (Continued)



(c) Fatigue crack morphology on specimen surface (growth direction from left to right)



9. PLASTIC ZONE OF A MIXED FATIGUE-CLEAVAGE CRACK IN SPECIMEN 3P-24 (Continued) FIGURE



FIGURE 9. PLASTIC ZONE OF A MIXED FATIGUE-CLEAVAGE CRACK IN SPECIMEN 3P-24 (Continued)



9. PLASTIC ZONE OF A MIXED FATIGUE-CLEAVAGE CRACK IN SPECIMEN 3F-24 (Continued) FIGURE





9. PLASTIC ZONE OF A MIXED FATIGUE CLEAVAGE CRACK IN SPECIMEN 3P-24 (Continued) FIGURE

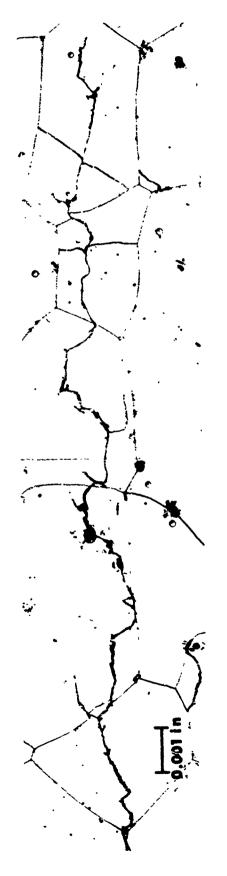


FIGURE 10. PLASTIC ZONE OF A MIXED FATIGUE-CLEAVAGE CRACK PRODUCED IN SPECIMEN 3P-25 BY 9930 CYCLES WITH A STRESS INTENSITY RANGE  $\frac{\Delta K}{\sigma_Y} = 0.7 \sqrt{1n}$  (CLEAVAGE PORTIONS ARE LABELED C)



(b) Plistic zone revealed in the interior on the specimen midsection

FIGURE 10. PLASTIC ZONE OF A MIXED FATIGUE-CLEAVAGE CRACK IN SPECIMEN 3P-25 (Continued)



(c) Morphology of the fatigue crack on the specimen midsection

FIGURE 10. PLASTIC ZONE OF A MIXED FATIGUE CLEAVAGE CRACK IN SPECIMEN 3P-25 (Continued)

FIGURE 10. PLASTIC ZONE OF A MIXED FATIGUE-CLEAVAGE CRACK IN SPECIMEN 31-25 (Continued)



COMPARISON OF FATIGUE CRACK- AND STATIONARY, VINGIN CRACK-NEAR-TIP PLASTIC ZONES: (a) tip of fatigue crack of Specimen 3P-25 growing at  $\sim 10^{-5}$  in, per cycle at  $\frac{10}{2}$   $\sim 0.5 \, \sqrt{1n}$ , and (b) virgin crack loaded to  $\frac{1}{2}$  = 0.4  $\sqrt{1n}$  and unloaded 41) FIGURE 11.

The light etching region near the crack tip is the region strained plastically in excess of 5%. This region has essentially the same shape in both cases including the small protuberance directly in front of the crack. Taking into account the difference in stress intensity levels, the 5%-strain-region surrounding the fatigue crack is still substantially larger as a result of the accumulation of plastic strain during reverse cycling. NOTE:

The crack branches seem to have an included angle  $\sim 90^\circ$  and are in all cases at least 10 x larger than the average crack growth per cycle, a sign that the branches represent not merely nucleation but many cycles of growth. The branches suggest there are two alternate and symmetric sites near the tip of the main fatigue crack for the cyclic growth, and that the two sites occasionally operate simultaneously. Similar observations have been reported and are discussed by Laird<sup>(33)</sup>. It is also interesting to note that some of the branches were arrested at grain boundaries (see Figure 9g).

Table 2 illustrates that the fatigue crack plastic zone size is essentially the same as that displayed by the slits that were loaded to the same peak stress intensity level. Previous experience shows that the size of the zone attending a sharp crack is essentially the same as the slit zone-size (when the extent of the zone is much larger than the slit-tip radius). Thus, the result derived from actual fatigue cracks and simulation experiment are the same: namely, there does not seem to be a significant difference in the plastic zone size of a cyclically loaded, growing crack and the monotonically loaded crack under plane strain conditions.

Figure 11 compares the near-tip strain distributions of a fatigue zone with a monotonic zone (41) at comparable stress levels. It can be seen that the light etching regions of both the monotonic zone display small protuberances directly in front of the crack. It seems likely that these are manifestations of the intensely strained region near the tip of the crack (see Figure 12a) recently treated by Rice and Johnson (29). The shape of light etching region near the tip can be rationalized in terms of the slip line field and the superposition of the two symmetric enclaves directly in front of the crack as shown in Figure 12b. The existence of protuberances in both zones is further evidence that the basic character of the strain distribution is not altered by cycling. The fatigue zone does display a larger light etching region and this is consistent with the superposition and accumulation of deformation from successive cycles.

#### **DISCUSSION**

The main finding here is there are no striking differences between fatigue crack zones—those produced in the foil and in the Fe-3Si plate by simulation and actual cracking—and the corresponding monotonic zones. While no radical differences were anticipated, the result is at odds with the expectation, discussed by Rice<sup>(28)</sup> that the steady-state cyclic zone is 1/4 the size of the stationary monotonic zone<sup>††</sup>. The discrepancy may simply be a consequence of the yield stress reduction affected by cycling (Bauschinger effect in the

The stress intensity level associated with the fatigue crack tip on the specimen midsection is less than the applied stress intensity level  $\left(\frac{\Delta K}{\sigma_Y} = 0.7\sqrt{\text{in}}\right)$  because the crack tunneled. The value quoted in Figure 12:  $\frac{\Delta K}{\sigma_Y} \sim 0.05\sqrt{\text{in}}$  is an estimate based on the actual size of zone, and  $\frac{\Delta K}{\sigma_Y}$  zone size relation developed for this DCB-specimen configuration (45).

Fe-3Si steel and cyclic softening in the cold worked foil). Cyclic stressstrain measurements for these materials are needed to establish whether the reductions are large enough to account for the zones.

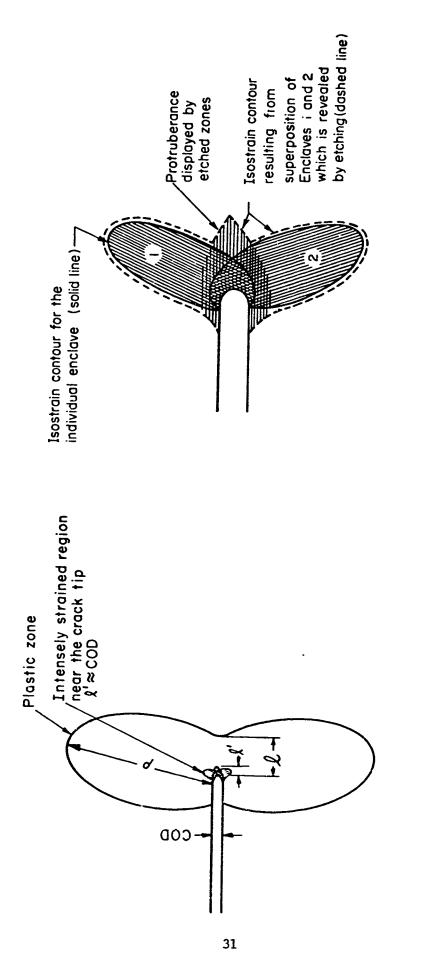
While questions remain, the observed similarities in size, shape and strain distribution do support the practice of applying the monotonic functional relations to the cyclic growth problem. For example, previous calculations and etching experiments indicate that the dimensions  $\rho$ ,  $\ell$ ,  $\ell^1$ , and COD under plane strain (see Figure 12a) can be approximated by the following expressions (28,29,41,46):

MONOTONIC ZONE	CYCLIC ZONE
$\rho \approx 0.2 \left(\frac{K}{Y}\right)^2  \dots  (1A)$	$\rho \sim 0.2 \left(\frac{\Delta K}{2Y}\right)^2 \qquad \dots  (1B)$
$\mathbf{\ell} \approx 0.03 \left(\frac{\mathbf{K}}{\mathbf{Y}}\right)^2 \dots$ (2A)	$\ell \sim 0.03 \left(\frac{\Delta K}{2Y_c}\right)^2 \dots (2B)$
L¹ ≈ COD	$L^{\dagger} \sim COD$ (3)
$COD \approx \frac{K^2}{2EY} \qquad \qquad (4A)$	$COD \sim \frac{\Delta K^2}{4EY_c} \qquad (4B)$

where K is the stress intensity, Y is the uniaxial yield stress, E the elastic modulus,  $\Delta K$  the stress intensity range,  $Y_c$  the cyclic yield stress and COD the crack opening displacement. The observation that  $\rho(\text{monotonic})\approx\rho(\text{cyclic})$  for the Fe-3Si Steel (see Table 2) suggests that Y  $\approx$  2Y $_c$ . As will be discussed later, this is only likely to be true in the lightly strained region that controls the outward reach of the plastic enclave, and not in the heavily deformed region that affects the value of  $\boldsymbol{\ell}'$  and COD. The number of cycles of plastic straining experienced by material in front of the crack is:

$$\Upsilon = \frac{\ell}{\frac{da}{dN}} \sim \frac{0.03\Delta K^2}{\frac{da}{dN} 4Y_c^2} \sim \frac{0.03\Delta K^2}{\frac{da}{dN} Y^2}$$
(5)

where N is the total number of cycles of plastic strain, and N the number of cycles of large plastic strains located within the intensely strained near tip region. (29) Estimates obtained in this way, assuming either  $Y_c = Y$  and  $Y_c = 2Y$ , are listed in Table 3 for 3 different materials. The relative invariance of the calculated values of N and N is an indication that crack growth rates are scaled to the dimensions of the plastic zone. A crude estimate of the total



SCHEMATIC OF MONOTONICALLY LCADED CRACK TIP PLASTIC ZONE: (a) plastic enclave identifying dimensions, (b) slip line field pattern and (c) near tip, high strain contours FIGURE 12.

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ESTIMATES OF \( \textstyle AND \( \textstyle \), THE NUMBERS OF PLASTIC STRAIN CYCLES AND \( \textstyle \), THE RATIO OF CRACK GROWTH PER CYCLE TO CRACK-OPENING DISPLACEMENT TABLE 3.

Material	Y (ksi)	Reference	da dN in/cycle	∆K ksi √in	Y <sub>c</sub> = Y	Ȣ	<b>a</b>	$\mathcal{N}^{x_c} = \frac{2x}{\mathcal{N}}$	۲, ک	<b>6</b> 0.
7075-T6	~ 72	Liu and lino(22)	10 -3 10 -4 10 -5 10 -6	37 24 11 4.6	10 (a) 30 70 120	~1(a) 2 4	1.0 0.5 0.25 0.1	5 (a) 15 35 60		>1 ~1 0.5 0.3
ну-130	~140	Barsom, et al. (24)	10 13 10 4 10 4 10 6	140 83 29 10	30 (a) 110 130 150	~1(a) 4 5 6	1.0 0.25 0.2 0.2	15 (a) 55 65 75	7000	>1 0 5 0 4 0 3
10N1 -Cr -Mo -Co	~190	Barsom, et al. (24)	10 -4 10 -5 10 -6	165 85 35 13	20 (a) 60 100 140	~ <sub>1</sub> (a) 3 5 7	00.3	10 (a) 30 50 70 100		0.7 0.3 0.3

(a) Probably an underestimate because plane strain conditions not satisfied.

accumulated plastic strain can also be derived:  $2\Delta \overline{\epsilon_p} \ \mathcal{N}' \sim 7^{\dagger}$ . This suggests that the flow stress for deformation near the crack tip could be raised substantially by strain hardening, and for this reason, estimates are also given in Table 3 for  $Y_c = 2Y$ .

Equations 1-4 are used in Appendix B to derive simplified expressions of the growth rates implied by the two mechanisms of crack advance that have currency:

(1) Irreversible Plastic Blunting (Kinematic Irreversibility) (32-34). The crack advances because the deformation attending the unloading part of the cycle is not an exact reversal of the deformation that blunts the crack during loading:

$$\frac{da}{dN} \sim \beta \frac{\Delta K}{4EY_C} \tag{7}$$

and 
$$\begin{cases} \frac{da}{dN} \sim 4 \left(\frac{\Delta K}{E}\right)^2 \\ \text{for } Y_c = 0.01E, \ \beta = 0.15 \end{cases}$$

(22,31) Damage Accumulation (22,31). The actual rupture of regions damaged by the succession of plastic strain cycles experienced by material in front of the crack:

DAMAGE ACCUMULATION

Large Strains:

Small Strains:

$$\frac{da}{dN} \sim \frac{\Delta \epsilon_{p}^{2}}{4B^{2}EY_{c}} \Delta K^{2}$$
 (8) 
$$\frac{da}{dN} \sim \frac{0.1}{B^{2}} \left(\frac{\Delta K}{E}\right)^{2}$$
 (9)

$$\begin{cases} \frac{da}{dN} \sim 25 \left( \frac{\Delta K}{E} \right)^2 & (8A) & \left( \frac{da}{dN} \sim 0.2 \left( \frac{\Delta K}{E} \right)^2 \\ \text{for } \frac{\Delta \overline{\epsilon}_p^2}{B_2} = 1, \ Y_c = 0.01E & \text{for } B^2 = 0.5 \end{cases}$$
 (9A)

<sup>†</sup>  $\bigcap$  '  $\sim$  5 for  $\frac{da}{dN}$  = 10<sup>-5</sup> and Y = Y<sub>c</sub>, and  $\Delta\varepsilon_p$ , the weighted average strain range close to the crack tip corresponds to the average strain in the monotonically loaded zone at distances smaller than the COD and is  $\sim 0.7(29)$ .

The quantity  $\beta \equiv \frac{da}{dN}/COD$  is the efficiency of the blunting process, B is a measure of ductility, and  $\Delta \overline{\epsilon}_p$ , the weighted average plastic strain range per cycle. These expressions are similar to the empirical growth rate law obeyed by a wide range of alloys in Regime No. 1 (see Appendix A and Tables 4 and A-1):

$$\frac{\mathrm{da}}{\mathrm{dN}} \approx \mathrm{A'} \left( \frac{\Delta \mathrm{K}}{\mathrm{E}} \right)^2 \tag{10}$$

where  $A' \sim 8$ . The experiments thus provide a way of examining the two mechanisms critically:

- a. Stress Intensity Dependence. Both mechanisms can account for the experimentally observed stress intensity exponent:  $m\approx 2$ , which applies to both the microscopic growth rate and the striation spacing in Regime No. 1. (See results of Broek(21) on 7075-T6 and Miller(26) on 4340 and 18 Ni-Maraging Steel in Table A-1.
- b. Absence of a Yield Strength Dependence. Tables 4 and A-1 illustrate that crack growth rates in Regime No. 1 are independent of the yield stress level for a wide range of materials. At first glance this result appears to favor the "small strain" damage accumulation argument since the other two treatments predict a dependence on Y<sub>C</sub> (Equations 7 and 8). However, Y<sub>C</sub> is not the ordinary yield stress but the cyclic yield or flow stress. Evidence is presented in Appendix C that the COD is most strongly influenced by the value of flow stress appropriate for the heavily strained material close to the crack tip. Since the rate of strain hardening varies inversely with strength level, it is possible that the cyclic flow stress at high strains approaches a common level, e.g., Y<sub>C</sub> ~ F for all materials. By way of this argument, which was suggested by Laird(47), the absence of a yield stress dependence can be rationalized with either mechanism (see Equations 7A and 8A).
- c. Modulus Dependence. Table 4 illustrates that when the resistance cyclic crack growth is normalized with respect to modulus in Regime 1, the  $\frac{\Delta K}{E}\Big|_{10}$  -5 values are nearly the same irrespective of alloy system. This feature, noted earlier by Bates and Clark is consistent with both mechanisms provided  $Y_c$  is a fixed fraction of E (see Equations 7A and 8A).
- d. Absence of a Ductility Effect. Table 4 shows that steels known to possess different levels of ductility still display the same resistance to crack growth. This seems to be inconsistent with the damage accumulation mechanism. However, it can also be argued here that the ductility in question is not the ordinary smooth bar value, but the value for a highly localized region under a complex state of stress, and this could be relatively invariant.

TABLE 4. SUMMARY OF SELECTED CYCLIC CRACK GROWTH RESISTANCE VALUES (See Table A-1 for Entire Compilation)

(Ks1) 54 93 95 95 140 163 191 218	3	ni 10 in	1	
54 93 95 140 140 at 163 F) Co 191 ng 218	(Kai Vin)	10-3 Vin	(Ksi Vin)	
93 95 140 140 F) Co 191 ng 218	~ 38	1,3	3	Present Study
95 140 140 F) Co 191 ng 218	39	1.3	1	Bates and Clark (15)
140 F) 163 Co 191 ng 218	33	1.1	1	Barson, et al, (24)
E) 163 (Co 191 218 218	29.5	1,0	92	
F) Co 191 ng 218	33	1.1	Ú8	Miller (18,26)
191 218				(23)
218	35	1.2	36	Parsom, et al.
,	31	1.0	80	
	252 36	1.2	80	Wei, et al. (2) and Carman
	č	•		and Katlin <sup>(3)</sup>
747 TTH	40	T• T	1 1	Carman and Katlin'
D6AC 241	32	1.1	~ 70	=
	,			(15)
5456-H321 37	12	1.2	ı	Bates and Clark, ""
2024-T3 50	9.5/13	1.0/1.3	29	Forman, et al, (%), Donaldson
				and Anderson (4), McEvily and
				Illg(2), Hudson and Hardvath(3),
2024-T3 50	10.5	1.1	16	Brock and Schiive(6)
	71	4.1		Brook (21)
	C.	0. [	20	=
(Argon)	14.5	2.5	•	Wet and Landes (27)
	2.5	90.00	22	
7079-16 65	12	1.2	¦ •	Bares and Clerk (15)
T1-6A1-4V	7.	o c	1	Ξ
	}	•	l	
!	į	•		(1)
70-30 Brass 18	14	6.0	•	McEvily, et al. "
70-30 Brass 92	14	6.0	ſ	=

- Efficiency of the Blunting and Damage Accumulation Mechanisms. Relatively little is known about the efficiency (degree of irreversibility) of the blunting process. A value of  $\beta \approx 0.15$  can be deduced directly in one case -- from photographs of a relatively short, highly strained crack in a polyethylene sheet published by McEvily and co-workers (49) --which define both the COD and  $\frac{da}{dN}$  This value is reasonably close to the  $\beta$ -values calculated for three alloys in Table 3 for  $Y_c = Y$ . The higher value, Y = 2Y, associated with the heavily strained hardened region, implies a greater efficiency, e.g.,  $\beta$  = 0.3-0.4. Such an efficiency may well be unrealistic, but this is not significant since the calculations of COD could well be in error by a factor of 2. Viewed in this light there is reasonably good agreement between the absolute prediction of the blunting mechanism and experiment (between Equations 7A and 10). By the same token, the predictions derived from the "large strain" damage accumulation treatment, Equation 8A, are also reasonable, although the values of  $\Delta \varepsilon_{\rm p}$  and B are at present very poorly defined. While the "small strain" damage accumulation treatment is least satisfactory for plane strain, it is roughly in accord with the results for the steel foil as pointed out in a preceding section.
- Transient Crack Growth Rates. The damage accumulation mechanism implies that the material in front of a crack is conditioned by a number of plastic strain cycles. Consequently, the crack advance per cycle at a given AK should increase gradually and only approach the steady state value after some number of stress cycles. Similar transients should be observed whenever AK is increased abruptly; a gradually decreasing rate when AK is abruptly reduced. This expectation is not confirmed by photographs of the striations produced by programmed loads that have been published by McMillan and Pelloux(50) and Christensen and Harmon(51) which show no signs of transients. This probably is the strongest argument against the damage accumulation mechanism. To the extent that strain accumulation from successive cycles contributes to the value of  $Y_c$ , transients having the opposite character should be associated with the blunting mechanism. While the striations mentioned above provide no clear-cut evidence of such transients either, the experiments of Hardrath (52) do show that a high pre-stress can retard the crack growth rate, and this is consistent with the expectation for blunting.
- g. Environmental Effects. The results of Barsom (23) on 12Ni-5Cr-3Mo and Wei and Landes (27) on 7075-T651 in Table A-1 show that certain "aggressive" environments greatly enhance the crack growth rate in Regime No. 1. This effect is difficult to understand in terms of the environments' effect on the flow stress and blunting and easier to rationalize in terms of rupture, and thus, damage accumulation.

Regime No. 1. The preceding paragraphs show that the main features of the low stress-righ cycle portion of the growth rate spectrum can be explained either in terms of the plastic blunting or the damage accumulation mechanism.

While neither mechanism is entirely satisfactory, irreversible blunting seems to be consistent with more observations at the present time.

Regime No. 2. The mechanism of growth in Regime No. 2, the high stresslow cycle range is not well defined. A possible clue is the observation that macroscopic (or average) crack advance per cycle exceeds the striation spacing(15.21,26). This can begin to occur when brittle particles or cher weak interfaces well in advance of the main crack front come under the influence of the plastic zone and rupture. These ruptures not only contribute to the crack advance per cycle, but may be sources of secondary fatigue cracks as suggested by the work of Erhardt and Grant (53). Thus the same mechanism operating in Regime No. 1 when augmented by an increasing length of active crack front (making A a stress dependent term) can account for (1) higher growth rates, (2) a greater stress dependence and (3) the discrepancy between growth rate and striation spacing (see Figure 13). Accordingly, AK1-2: the stress intensity range at the transition from Regime No. 1 to Regime No. 2 should depend on the spacing and properties of specific microstructural features which may vary from alloy to alloy and heat to heat. Consistent with this, the AK1\_2 values listed in Tables 4 and A-1 are not constant and do not seem to correlate with yield strength. At higher stress intensity levels, the approach to general yielding and transition from plane strain to plane stress may also accelerate the crack growth rate.

This admittedly tentative picture could have important implications with respect to the metallurgical factors and alloy development. To the extent that environmental effects are not involved, neither Equation 10 nor the compilation of AK' values shows any systematic effects of composition, microstructure, or the usual thermal-mechanical treatments. It is true that metallurgical contributions amounting to  $\sim$   $\pm$  20% of  $\Delta$ K $_{10}$ -5 (corresponding with  $\sim$  50% change in total life) could be obscured by the scatter of the data. Beyond this, the prospects of improving crack growth resistance in Regime No. 1 by the usual techniques do not seem bright. It is possible \_nat large improvements could be realized by metallurgical changes that are designed to eliminate susceptibility to aggressive environments. Special microstructures, such as composites with interfaces normal to the crack front that delaminate preterentially, may offer advantages. The micrographs in this report contain some evidence that fatigue cracks are impeded by grain boundaries, but it is not clear whether enough boundaries can be introduced to affect the average crack growth increment significantly (54).

In contrast, the behavior in Regime No. 2, especially the value of  $\Delta K_{1-2}$ , could be quite sensitive to the microstructure, to the strength and spacing of hard particles and to other metallurgical factors if the explanation offered here is correct. Furthermore, significant improvements in total life can be achieved simply by increasing  $\Delta K_{1-2}$  and postponing the higher growth rates generated in Regime No. 2. To test this possibility, it would be useful to compare  $\Delta K_{1-2}$ -values obtained for a single phase alloy with modifications involving different distributions and morphologies of hard particles.

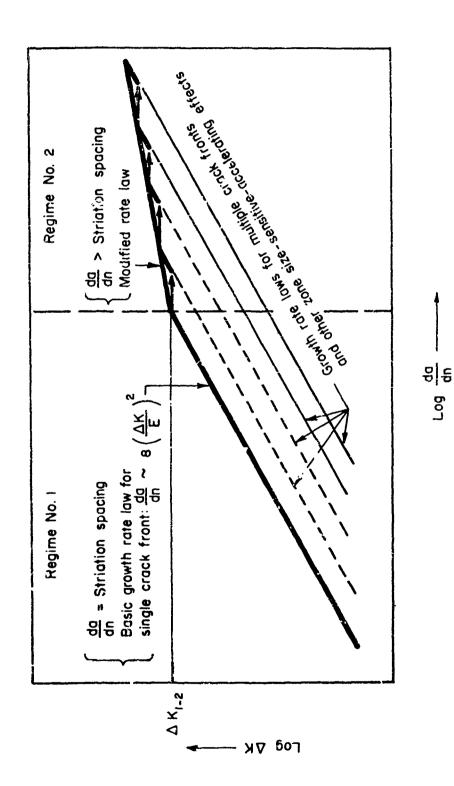


FIGURE 13. SIMPLIFIED SCHEMATIC REPRESENTATION OF THE ORIGIN OF REGIMES NO. 1 AND NO. 2

#### CONCLUSIONS

- 1. The interferometric technique applied to the steel foil and the etching technique for Fe-3Si can provide useful insights to the size of the fatigue crack plastic zone and the strain distribution within it. Preliminary results show that the cyclic loaded zone in the foil (plane stress) and in the Fe-3Si plate (plane strain) are essentially the same size as the corresponding "monotonic" zones. This result suggests that the cyclic flow stress at small strains is roughly 1/2 of the ordinary tensile yield stress consistent with the Bauschinger effect.
- 2. The observed similarities in the size, shape and strain distribution within the cyclic and monotonic zones lend support to the practice of adapting functional relations for the monotonic case to the fatigue problem. Estimates of the monotonic zone size and crack-tip displacement indicate that the material in front of a fatigue crack is subjected to about 50-150 cycles of reverse plastic deformation, of which about 2-7 cycles involve the large plastic strains experienced within the intensely deformed neartip region. The crack advance per cycle is estimated to be 10-50% of the crack opening displacement.
- 3. The compilation of cyclic crack growth-rate measurements for steels, aluminum alloys and a titanium alloy provides evidence for 2 growth-rate regimes. Regime No. 1 involves the low stress intensity values and low growth rates usually  $< 10^{-4}$  in. per cycle which correspond with the striation spacing, and a stress exponent  $m \approx 2$ . Regime No. 2 is the high stress-high growth rate portion of the spectrum usually  $> 10^{-5}$  in. per cycle and a larger stress exponent,  $m \gtrsim 4$ .
- 4. The cyclic crack growth resistance in Regime No. 1 is roughly proportional to the elastic modulus but relatively insensitive to yield strength, composition and ductility. Growth rate measurements on a variety of steels, aluminum alloys and Ti-6Al-4V can be approximated by the expression  $\frac{da}{dN} \approx 8 \left(\frac{\Delta K}{E}\right)^2, \text{ where E is the modulus. While neither the blunting or damage accumulation mechanisms easily account for all of the features of Regime No. 1, more of the existing observations can be rationalized by blunting .$
- 5. Some of the features of Regime No. 2 can be explained by the mechanism responsible for Regime No. 1 together with the idea that crack growth is further enhanced by an increasing length of active crack front. This is thought to occur when the plastic zone is large enough to produce cracks at hard particles and other weak interfaces ahead of the main crack front which then also extend by cyclic growth and connect with the main crack. Accordingly, improvements in fatigue life could be realized by postponing the transition from Regime No. 1 to No. 2, and this may be related to metallurgical factors.

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# APPENDIX A -- Compilation of Cyclic Crack-Growth Rate Measurements

Noticeable departures from a simple exponential crack growth relation  $\frac{da}{dN} = A(\Delta K)^m$  were noted by  $\text{Liu}^{(55)}$  as early as 1964. In the interim, most workers have tended to gloss over irregularities and have attributed a single mechanism,  $m \sim 4$  to the entire growth rate spectrum(56). However, the recent work of Bates and Clark(15), Barsom, et al.(24) and measurements of the stress dependence of the striation spacing(15,21,26) offer good evidence for two regimes of the behavior:

Regime No. 1. This is the low stress-low growth rate portion of the spectrum, e.g., usually less than  $10^{-4}$  in. per cycle, where the macroscopic growth rate coincides with the striation spacing (15,21,26). As shown in Tables A-1 and 4, growth rates for a wide range of alloys are approximately described by the following empirical expression:

$$\frac{\mathrm{d}\mathbf{a}}{\mathrm{d}\mathbf{N}} \approx \mathbf{A}^{\dagger} \left( \frac{\Delta \mathbf{K}}{\mathbf{E}} \right)^{2} \tag{A-1}$$

where  $A' \sim 8$ .

Regime No. 2. This is the high stress-high growth rate position of the spectrum, e.g., usually more than  $10^{-5}$  in. per cycle. In this range growth rates tend to exceed the striation spacing and displays a greater stress sensitivity, e.g.,  $m \ge 4$ .

These two regimes are shown schematically in Figure A-1 which also defines the symbols used here to characterize the two regimes and the transition. The compilation of crack growth rate measurements in Table A-1 represents a conscious effort to identify two types of behavior and to assign them to Regimes No. 1 and No. 2 on the basis of the exponent. In some cases, neither a clear-cut change in behavior nor a characteristic m could be detected, and the Regime assignment is arbitrary. The compilation is mainly for growth rate measurements in air, with R  $\sim$  0, but Barsom's  $^{(23)}$  measurements for 12Ni-5Cr-3Mo Steel in NaCl and Wei and Landes  $^{(27)}$  results for 7075-T651 in Argon and  $_{20}$  are included to illustrate the potency of "aggressive" environments.

It should be noted that when the coefficient A is evaluated empirically, it involves the exponent m:  $A = \begin{pmatrix} \frac{da}{dN} \\ i \end{pmatrix}_{i} \Delta K^{-m} \text{ (where } \begin{pmatrix} \frac{da}{dN} \\ i \end{pmatrix}_{i} \text{ is the growth rate corresponding to } \Delta K_{i} \text{).} Since m is frequently not known with precision, uncertainties in m are magnified in A even where <math>\Delta K_{i}$  is established accurately. For this reason  $\Delta K_{i}$  is a more reliable index of performance than A, and  $\Delta K_{i}$  and m (rather than A and m) are used here to characterize the growth rates and for purposes of comparison.

DATA
GROWTH
CRACK
CYCLIC
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SUPPLARY
A-1.
TABLE

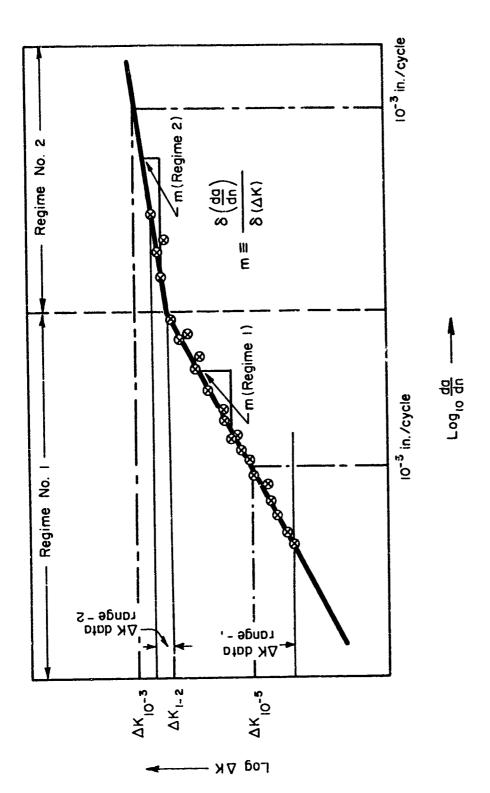
					Transition	- Tuota	Regi	Regime No, 2		
Materie 1	Y (K81)	OK-DATA RANGE	SE CK10-5	Œ	0K <sub>1-2</sub> (K•1√In)	* (E)	OK-DATA RANGE (Keivīr)	<sup>∆K</sup> 10-3 (Ks1√In)	8	References
		THE STREET	7110 10 201							
				Part 1,	1, Steels					
Fe-3S1/_	25	1	38	•	•		•	• }	• ]	Present Study (19)
A-302B(C)	59	•	1	•	• :	• ;	40-120		1.1	Mobray, et al
A-302B(G)	21	•	•	•	0 <del>7</del> ~	8.0.	011-07		7.0	: 3
A-508-2(d)	60	• 1			• 1	• :	60-10	`` } } ``	~ ~ ~ ~ ~ 3.5	
A-508-2	73	10_100(e)	(e) (g) (g) (R)	2.1/2.2(e)			2			Bates and Clark (15)
3.X X.	) (T	30-150	34		•	•	•		•	
HV-80 (a)	, r 00	21-61	38		1		•	•		Barsom, et al (22)
H72-80(b)	, 6 	23-50	33			•	•	•	•	(11)
HY-100.Steel	109	5.6-26	20/23		1	•	•	•	•	Yokobori, et 41.
HY -130(4)	139	19-68	28						• 5	Barsom, et al (**)
HY-130 (b)	140	4292	29.5			0.66	92-150	140	8/9	(35)
4340 (ANNEALED)	09	14-20/25	20			0.3/0.4	20/25-35		4.0 (a) (e)	11u and 11p9
4340 (Q & T @ 1400F)	69	• 6	55.			•	25-70		4/3.6	, , 197111
4340 (Q & T @ 1000F)	163	20-80	7	7.4/2.20	2 G 2 G	94.0	017-06			=
4340 (4 & I @ 300F)	177					3	20-75	~ 20	6.73	
100 A G C C C C C C C C C C C C C C C C C C	180		•	•	•	. 1	40-100		3,7	Crooker, et al As
977 TOO 1776	175/180	08-57	~ 34	~ 1.85	•	•	•		•	Bates and Claris, 112,
10N1 -Cr -Mo-Co (a)	191	50	35	2,25	86	0.45	86-140		3.8	Barsom, et al('4)
10N1 -Cr-Mo-Co (b)	182	30≖60	30	2.24		•	•	•	•	* :
12N1-5Cr-3Mo(E)	185	21-60	31	2.16	•	•	•	•		: 1
12N1-5Cr-3Mo(b)	184	25-60	28	2.18		•	•		• ;	(14)
12N1-5Cr-3Mo	194	33-80	33	2.3	80		80-105	}	3.	Schwab (22)
12N1-5Cr-3Mo (NAC1, 600CPM)	184	15-53	26	~ 2.0		1	•		•	Barson
12N1-5Cr-3Mo (NaC1, 60CPM)	184	18-55	23	, Z.0		•	•		•	
12N1-5Cr-3Mo (NACI, 6CPM)	184	20-47	91	~ 2.0 				•	•	7 10 10 10 10 10 10 10 10 10 10 10 10 10
12N1-Mareging	180	25-100		2.3		•		•		crooker and Lange (16)
18N. Western	180	06-67	2	2.3		,	•	•	•	= ;
18N1-Maraging	218	20-80	3 5	2.3/2.5/1.42(e)	~ 80		80-100	~ 110	•	Miller (18, 25)
18N1 -Maraging	245	20-100	35	7, 88/1.74 (e,			•	~ 150		-
18N1 -Maraging	246/252			2.4	08	0,32	80~100	•	•	Wei, et al " and
										Carman and Katiin (8)
18Ni -Marasine	295/308	18-60	35	2.6	•	•	•	•	•	'let, et al(12) and
			<b>:</b>	•						Carmen and
н 11	208	•		•	•		21-29		6.62	Kat 11n ( ) M411er (18, 26)
H II	242		35				27-60	~ 120 ~	4.0	Carman and
	,	,								Katling
D6AC	241	32-80	32	~ 2.8	. 70	•	ŧ	•	•	Carmen and Katiin(8)
(a) 1-inthick plate (b) 2-inthick plate (c) Thickness direction (d) Axial direction										
	cing									

TABLE A-1. SUMMARY OF CYCLIC CRACK GROWTH DATA (CONTINUED)

					Two				
		Regime	Regime No. 1		1	1		:	
Material	Y (Ks1)	OK-DATA RANGE (Ksivin)	ΔK <sub>10</sub> -5 (Ks (γ(i))	E	∆K <sub>1-2</sub> (Ks1√in)		AK-DATA RANGE	NGE OK10-3	m References
							(III AT SW)	(KSIVin)	
				Part 2.	Part 2. Aluminum Allovs	lovs			
7075-T6	02	5-2.2	7.5	2,3	22	11			
7075-T6	70	6-24	01	2,4	24	·	75-77		S Forman, et al and McEvily and Illa(2)
7075-I6 (R = C,3)	0/	5-18	œ	2.8	; <u>«</u>	t %	59 <del>-</del> 57	40 ~ 5.7	
7075-T6	75,5	6-13	•	2.75	? :	67.0	18-22		Forman, et al(9) and Brock and Schaus(6)
	0 0 0	5-20	10	2/2(0)	5 S		13-30 20-28	~ 22 ~ 5	Hudson and Scarding (25)
7075-T651 (H <sub>2</sub> O)	2,2	6-18	7. 8 0. 9	9,0	•	•			Brock to
	20	7-29	9.5/13	2,6/3	29	85 0	• 6		
						3	0/-67	8.9 ~ 6.8	ă
* *	20	5-2%(a)	10 5 (8)	!	(*)				Hudien and Hardrath(5)
2024-T3 (R = 0.3)	20	4-10	8 8	3.9	22 (4) 18	0.44(4)	$22^{(a)}$ - $25^{(a)}$	$\sim \frac{33^{(a)}}{23}$	Merbull Illg and McEvilv(3)
2024-T3 (R = 0,25)	ç	4.11	•	,	Ì			- 57	Forman, et al(9)
	20	4-1:		6.2	16	0,32	16-40	35 ~ 4.5	brock and Schijve(9)  Brock and Schijve(6)
7079-T6	20	5-26	14	2.5/2(b)		0.24	12-20		Liu and Jino(22)
7079-T6	0 00	8-30	12	2.6/2.1(6)			cc-27	2. ~	Brock, 2D
7079-T6	0,	15-50		~ 2.7	ដ	0.19	13-20	$\sim 20 \sim 7$	Swanson Ar 11(10)
125H-9546	37	9-20	12	2,1/2,1		• (	•		Smith, et a1 (20)
						,	•		Bates and $Clark^{(15)}$
				Part 3,	Titanium Alloys	870			
T1-6A1-4V T1-6A1-4V	127 121	13-35 40-66	15 38	3/2(b) ~ 4	l t	1 5	• •	• •	Bates and Clark(15) Hall(17)
				Part 4. M	Miscellaneous Alloys	Alloys			
70-30 Brass (Annealed) 70-30 Brass (cold worked)	18 92		14 14	N. N.					MCEVILY, et al (7)
									=

<sup>(</sup>a) 'Value3 quoted represent 1/2 the value reported because R = 1

<sup>(</sup>b) Based on striation spacing



SCHEMATIC OF STRESS-INTENSITY-GROWTH RATE SPECTRUM WITH 2 REGIMES AND DEFINITIONS OF TERMS IN TABLE A-1 FIGURE A-1.

# APPENDIX B -- Simplified Formulations of the Cyclic Crack Growth Mechanisms

A. Irreversible Plastic Blunting. This process can be formulated approximately by noting that  $\frac{da}{dN}$ , the crack advance per cycle will be a fraction  $\beta$  of the OD (Equation 4B):

$$\frac{da}{dN} = \beta \cdot COD \approx \beta \cdot \frac{\Delta K^2}{4EY_C}$$
 (B-1)

where  $\beta$ , the efficiency (or degree of irreversibility) of the blunting process is regarded as constant to a first approximation.

B. <u>Damage Accumulation</u>. This mechanism is described in essentially the same way as proposed by Lehr and Liu(31) by combining Equations 5 and 6 with the Manson-Coffin fatigue damaged law(35,36,47):

where  $\Delta \overline{\varepsilon}_p$  is regarded as the weighted average plastic strain range and B the fracture strain for monotonic loading: This gives two results:

small strain 
$$\frac{da}{dN} \approx \frac{0.03\Delta \overline{\epsilon}_{P} \Delta K^{2}}{4B^{2}Y_{c}^{2}} \sim \frac{0.1}{B^{2}} \left(\frac{\Delta K}{E}\right)^{2}$$
 (B-3)

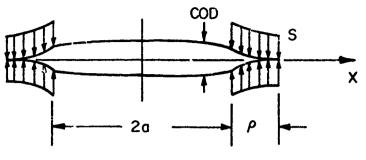
large strain 
$$\frac{da}{dN} \approx \frac{\frac{\Delta \bar{e}_{P}^{2} \Delta K^{2}}{4B^{2}EY_{C}}}{(B-4)}$$

In the first,  $\Delta \overline{\varepsilon}_p^{-2}$  is the strain range averaged over all cycles, which and Liu estimate as  $\Delta \overline{\varepsilon}_p \approx 4 \, \frac{Y_c}{E}$  and is consequently a relatively small strain, e.g.,  $\Delta \overline{\varepsilon}_p \approx 0.02$ . In the second, only the number of cycles involving large strains is considered and  $\Delta \overline{\varepsilon}_p \sim 0.5^{(29)}$ .

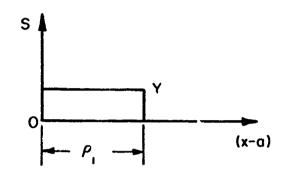
# APPENDIX C--Augmented Dugdale Model

At the tip of a fatigue crack there exists a very heavily strained region where  $Y_c$ , the cyclic yield stress, may be greatly in excess of Y, the yield stress for ordinary monotonic loading. Simultaneously,  $Y_c$  may be less than Y in the remaining, lightly strained region of the plastic zone as a result of the Bauschinger effect. A model of this behavior is given in Figure C-1. Figure C-la shows the Dugdale Model with the tips of the crack being pressed together by a traction denoted S. The region -a<x<a represents the crack, a<x<a + \rho\$ the plastic zone. Usually S is set equal to a constant yield stress as in Figure C-lb. For the case of the fatigue crack, the distribution shown in Figure C-lc is post

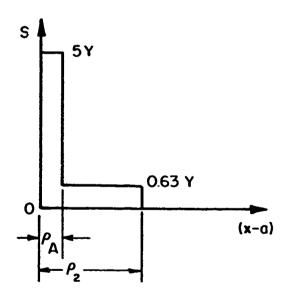
The two Y<sub>C</sub>-levels represent the small heavily strained region where Y<sub>C</sub> = 5Y and the lightly strained region where Y<sub>C</sub> = 0.63Y. According to the calculation, which was made by methods outlined in References (57) and (58), COD (Case 2) = 1/5 COD (Case 1), when  $\rho_1 = \rho_2$ ,  $\rho_2/\rho_A = 5.5$  and  $\frac{T}{Y} = 0.25$ , where T is the nominal applied stress. Thus, there is a large difference in COD even though the plastic zones are the same size. Also the COD for the fatigue crack model is comparable to a much stronger material even though the heavily strained region occupies a small part of the total plastic zone. The model illustrates that the COD is strongly influenced by the value of Y appropriate for material closest to the crack tip.



a. Dugdale Model



b. Case I - Constant Flow Stress



c. Case 2-Model of Flow Stress Distribution for Fatigue Crack

FIGURE C-1. AUGMENTED DUGDALE CRACK MODEL

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3. ABSTRACTThis report describes experimen			
atigue crack and its relation to the zon compilation of cyclic crack growth meas of the zone studies and the data compilate of cyclic growth and the contribution of Two techniques, etch pitting and intrones produced by growing fatigue cracks worked steel (under plane stress), respectively crack growth under controlled conduittent load reversals is demonstrated. deformation generated by each loading cyclic crack growth under controlled conduities the conduction generated by each loading cyclic crack growth under controlled conduities the conduction generated by each loading cyclic crack growth under controlled conductions.	urements is als tion are examine metallurgical for the cerferometry, are in Fe-3Si Steel ctively. In additions by cutti	o present d with re factors. e used to (under pliticn, a ing-in sli results	ed. Preliminary results espect to the mechanism of reveal the plastic plane strain) and in columethod of simulating its under load with inte indicate that the plast
The results provide a basis for extended crack to the fatigue crack problem of plastic strain cycles experienced by the sand also the crack-tip displacement. The treatment of cyclic crack growth which is their analysis is in good accord with measures attribute striction spacing observed in Registration.	n. Estimates and the material in the values obtained based on Mansonsurements of the contract	e made in front of d lend su on-Coffin he cyclic	n this way of the number a growing fatigue crack apport to Lehr and Liu's type damage accumulatio cracking resistance and

Both the theoretical treatment and the bulk of the test data show that the resistance to fetigle crack growth in Regime No. 1 is relatively insensitive to composition and metallicital structure, provided environmental effects are not involved. In contrast, significant improvements in fatigue life could be achieved simply by postponing the onset of Regime No. 2, a feature that should (Unclassified) depend on metallurgical variables.

Security Classification

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Unclassified
Security Classification